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# Understanding the Capability of Future Direct-imaging Observations to Quantify Atmospheric Chemical Effects of Stellar Proton Events

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## Abstract

Models developed for Earth are often applied in exoplanet contexts. Validation in extraterrestrial settings can provide an important test of model realism and increase our confidence in model predictions. NASA's upcoming space-based IROUV telescope will provide unprecedented opportunities to perform such tests. Here, we use the Planetary Spectrum Generator to simulate IROUV reflected-light spectroscopic observations of flare-driven photochemical changes produced by the Whole Atmosphere Community Climate Model, part of the Community Earth System Model framework. We find that  $NO<sub>2</sub>$  is the most observable gas to target, and integrating the signal for two days following the flare and comparing to a baseline of preflare data would achieve the highest signal-tonoise ratio. The  $NO<sub>2</sub>$  response is much larger for K-star tidally locked planets than G-star rapidly rotating planets and does not depend strongly on  $O_2$  level. The NO<sub>2</sub> response should be observable for planets within  $3-4$  pc independent of the phase angle since the amount of reflected light is larger at smaller phases, but the  $NO<sub>2</sub>$ concentration is low near the substellar point. This work outlines a methodology for validating and ground-truthing atmospheric chemistry models developed for Earth that could be useful for the numerical exploration of exoplanets.

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#### 1. Introduction

Atmospheric chemistry and photochemistry are important factors in understanding exoplanetary habitability and performing atmospheric retrievals. Photochemical modeling of rocky exoplanets was first performed in one-dimensional (1D) models (e.g., Segura et al. [2005;](#page-5-0) Kaltenegger & Sasselov [2009](#page-5-0); Segura et al. [2010](#page-5-0); Hu et al. [2012;](#page-5-0) Grenfell et al. [2013](#page-4-0); Gao et al. [2015](#page-4-0); Kozakis et al. [2018](#page-5-0)) that provided key insights into the effects of stellar spectral energy distribution, stellar variability, and stellar flares. More recently, such modeling efforts have been expanded to the 2D and 3D regimes, usually with atmospheric chemistry and photochemistry as additional subcomponents in global Earth system or general circulation models (Chen et al. [2018,](#page-4-0) [2019,](#page-4-0) [2021](#page-4-0); Braam et al. [2022;](#page-4-0) Cooke et al. [2022](#page-4-0); Ridgway et al. [2023](#page-5-0); Tsai et al. [2022](#page-5-0)). While the efficiency of single-column models allows one to explore a large parameter space of planetary characteristics, higher dimensional (2D/3D) models are able to simulate the complex interplay among atmospheric dynamics, clouds, radiation, and chemistry and hence provide results with enhanced realism.

Upcoming direct-imaging missions of potentially habitable worlds provide an unprecedented opportunity to test model realism and predictions of terrestrial photochemical models in

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extrasolar settings. Evaluating the effects of stellar flares on modeled atmospheric chemistry against observational measurements is one such promising potential method. For instance, Chen et al.  $(2021)$  $(2021)$  $(2021)$  studied the effect of stellar flares on the atmospheric chemistry of terrestrial planets using the National Center for Atmospheric Research Whole Atmosphere Community Climate Model (WACCM), a well-studied and wellvalidated terrestrial Earth system model. They found that stellar flares can affect concentrations of  $NO_x$ ,  $HO_x$ , and  $O_3$ , but did not simulate reflected-light observations to determine whether these changes are potentially observable.

The Planetary Spectrum Generator (PSG; Villanueva et al. [2018](#page-5-0)) is a radiative transfer model that can can simulate spectral observations by current and future telescopes. PSG has been used to simulate both transit spectroscopy (Fauchez et al. [2019;](#page-4-0) Komacek et al. [2020](#page-5-0); Suissa & Wolf [2020;](#page-5-0) Haqq-Misra et al. [2022](#page-4-0)) and reflected-light spectroscopy (Checlair et al. [2021;](#page-4-0) Kopparapu et al. [2021](#page-5-0)) for terrestrial planets. In the case of reflected-light spectroscopy, PSG parameters can be chosen to simulate a space telescope that may not launch for decades. PSG provides us with the capability to determine whether the changes in atmospheric chemistry due to stellar flares are potentially observable.

In this paper, we use PSG to simulate reflected-light spectral observations of the response of the atmospheric chemistry of a terrestrial planet to stellar flares as simulated by Chen et al. ([2021](#page-4-0)). We assume a 6 m LUVOIR-like telescope as a proxy for NASA's IROUV telescope that will be capable of directly imaging Earth-like planets (National Research Council [2021](#page-5-0)). We find that  $NO<sub>2</sub>$  yields the largest signal-to-noise ratios of any

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simulated molecule, and that changes in  $NO<sub>2</sub>$  could be observable for planets within 3–4 pc. Detections would be possible for planets orbiting K stars, regardless of their background oxygen concentration, but not for planets orbiting G stars. The K-star planets simulated by Chen et al. ([2021](#page-4-0)) are tidally locked and they found that their  $NO<sub>2</sub>$  concentration is lower near the substellar point than the antistellar point. As a result, we find that the signal-to-noise ratio  $(S/N)$  is relatively insensitive to phase. This paper is organized as follows: in Section 2 we describe our methodology and in Section 3 we give our results. We discuss our results in Section [4](#page-4-0) and conclude in Section [5.](#page-4-0)

#### 2. Methodology

We examine the observational implications of flare-induced chemistry on high-mean molecular weight atmospheres by postprocessing the results of Chen et al. ([2021](#page-4-0)), who simulated  $O_2$ -rich and  $O_2$ -poor atmospheres for planets orbiting both G and K stars. We do not analyze the M-star cases because these planets are likely to fall within the inner working angle of a direct-imaging telescope. The planets have a modern Earth continental configuration. The G-star cases have a modern Earth rotational period (24 hr) and orbital period (1 yr). The K-star cases are tidally locked in a 1:1 spin:orbit state (rotational and orbital period of 92 Earth days) with the substellar point at the equator and 180° longitude (centered on the Pacific Ocean). The atmospheric surface pressure is 1 bar and we consider two levels of atmospheric oxygenation: (1)  $O_2$ -rich (21%, modern Earth-like) and (2)  $O_2$ -poor (1%, proterozoic Earth-like). WACCM calculates atmospheric chemistry using the Modules for Ozone and Related Chemical Tracers chemical transport model (Kinnison et al. [2007](#page-5-0)), which has a network of 217 reactions. The horizontal resolution is  $1^\circ.8 \times 2^\circ.5$  and there are 66 vertical levels with a vertical resolution of 0.5–2 km below the stratopause and roughly half a scale height above the stratopause.

We use the PSG (Villanueva et al. [2018](#page-5-0)) with the Global Exoplanet Spectra add-on to simulate reflected-light spectroscopic observations of these WACCM simulations. We choose parameters appropriate for the direct-imaging IROUV space telescope proposed in the NASA Decadal Survey (National Research Council [2021](#page-5-0)), which will directly image Earth-like exoplanets. To approximate IROUV, we use the LUVOIR PSG option with a diameter of 6 m. We use the PSG spatial binning option of 3, corresponding to  $3 \times 3$  (lat x lon) binning. A spatial binning of 3 is the smallest option available in PSG. Our results are not sensitive to this choice. For example, we find that the S/N is only  $\approx 5\%$ lower in a test where we used the maximum binning, so that all spatial variability is averaged out. We use the spectral parameters of the Sun for the G-star case and the spectral parameters of a K6V for K-star case. PSG allows us to specify the phase of the simulated observation, which we vary below.

We calculate the  $S/N$  for the retrieval of a molecule in the following way. First we use PSG to calculate the signal and noise, in terms of spectral intensity, in two cases: (case 1) including all molecules and (case 2) including all molecules except the molecule of interest. We then define the signal for the molecule of interest,  $S_i$ , as the difference between the spectrum in case  $2$  and case  $1$ , where the index  $i$  refers to wavelength bin. Each signal is associated with a noise,  $N_i$ , produced by PSG in case 2 (the noise for case 1 is almost identical). We then integrate the signal and noise over a number of days before the flare ( $S_{0i}$  and  $N_{0i}$ ) and after the flare ( $S_{fi}$  and  $N_f$ ). To integrate the signal we take its algebraic mean, while for the noise we use the formula  $\frac{1}{N^2} = \sum_{k}^{n} \frac{1}{N_k^2}$  where k is the day. We then calculate the  $S/N$  for the measurement in a given wavelength bin  $(\sigma_i)$  as

$$
\sigma_i = \frac{S_{fi} - S_{0i}}{\sqrt{N_{fi}^2 + N_{0i}^2}}.
$$
\n(1)

Finally, we combine all wavelength bins using the following formula  $\sigma = (\sum_i (\sigma_i)^2)^{\frac{1}{2}}$ .

We define a molecule with detectable variability as a molecule whose change in abundance in response to the flare can be detected with an  $S/N$  of at least 3. We checked  $H_2O$ ,  $CO<sub>2</sub>$ ,  $CH<sub>4</sub>$ ,  $O<sub>3</sub>$ ,  $NO<sub>2</sub>$ ,  $N<sub>2</sub>O$ , and  $H<sub>2</sub>$  and found that  $NO<sub>2</sub>$  is the only molecule with detectable variability among these. We also note that molecules  $HNO<sub>3</sub>$ ,  $NO+$ ,  $HO<sub>2</sub>NO$ ,  $H<sub>2</sub>O<sub>2</sub>$ , and  $NO<sub>3</sub>$  are produced by WACCM, but are not fully supported by PSG, so we were not able to investigate their detectability.

To validate our methodology, we reproduce the  $NO_2 S/N$ versus integration time generated by PSG in 1D configuration shown in Figure 6(b) of Kopparapu et al. ([2021](#page-5-0)). Like Kopparapu et al. ([2021](#page-5-0)), we use LUVOIR with a 15 m diameter for this calculation. Kopparapu et al. ([2021](#page-5-0)) use a modern Earth-like planet, so we use the  $O_2$ -rich G-star case from the Chen et al.  $(2021)$  $(2021)$  $(2021)$  simulations. The NO<sub>2</sub> concentration in the two models is similar and we calculate signal-to-noise ratios  $(S/Ns)$  that are 10% smaller than Kopparapu et al. ([2021](#page-5-0)) for a given integration time (Figure [1](#page-2-0)). Given that there are many PSG parameters that Kopparapu et al. ([2021](#page-5-0)) may have chosen to be slightly different from our calculation, we consider this an acceptable validation of our methodology.

## 3. Results

We do not find detectable variations in  $NO<sub>2</sub>$  for the planets orbiting G stars, so we focus on the tidally locked K-star planets in what follows. Figure [2](#page-2-0) shows the proton fluence (solar flare activity),  $NO<sub>2</sub>$  column, and simulated  $NO<sub>2</sub>$  S/N for an Earth-sized, oxygen-rich planet orbiting a K-star that is 2 pc from Earth and observed at a phase of 45°. Large responses in NO2 follow many flares with the largest of them being at days 58, 171 and 280, all of which lead to detectable signals. Some flares do not cause a large  $NO<sub>2</sub>$  enhancement and/or lead to small  $S/Ns$ . This is because  $NO<sub>2</sub>$  production depends on the presence of other photochemically produced species, such as NO and  $O_3$ , which are not always available. Additionally, there are some S/N peaks when we include clouds in the calculation that do not appear to be caused by a flare or increase in column NO2 abundance, with the most prominent examples occurring on days 85, 107, and 225. We believe this is related to an error that the PUMAS radiation scattering scheme used by PSG returns, "maximum asymmetry has been capped." On days where there is a flare and a resulting  $S/N$ , the magnitude of the  $S/N$  is similar with and without clouds (Figures [2](#page-2-0) and [3](#page-3-0)), so this issue should not substantially affect our main results.

We next discuss the optimal observing strategy for these detectable  $NO<sub>2</sub>$  responses to flares. One important consideration is the amount of expensive telescope time necessary to make an observation. Figure  $3$  shows the  $S/N$  for detecting

<span id="page-2-0"></span>

Figure 1. NO<sub>2</sub> vertical profile from Kopparapu et al. ([2021](#page-5-0)) and our O<sub>2</sub>-rich G-star planet (left). Signal-to-noise ratio (S/N) for a detection of NO<sub>2</sub> for both models assuming a LOIVOIR-type telescope with a 15 m diameter and only considering UV data, as in Kopparapu et al. ([2021;](#page-5-0) right).



Figure 2. Time series of the proton fluence (stellar flares, top), global-mean NO<sub>2</sub> column (middle,black), NO<sub>2</sub> column at the antistellar point (middle, red), NO<sub>2</sub> column at the substellar point (middle, blue), and  $S/N$  of  $NO<sub>2</sub>$  including clouds in the calculation (bottom, blue) and excluding clouds (bottom, black) for an Earthsized, oxygen-rich planet orbiting a K star that is 2 pc from Earth and observed at a phase of 45°.

changes in  $NO<sub>2</sub>$  as a function of integration period after the flare for two integration periods before the flare (5 and 10 days). The integration period before the flare is necessary to

establish a baseline  $NO<sub>2</sub>$  value, and our results are not very sensitive to this integration time. We find that the optimal integration period after the flare is two days. This results from

<span id="page-3-0"></span>

Figure 3. NO<sub>[2](#page-2-0)</sub> S/N for the flare at days 58 in Figure 2 for different integration periods. We integrate the signal before the flare over 5 (dashed line) or 10 (solid line) days and after the flare has begun over a range of days (horizontal axis). We perform calculations both including clouds (blue) and without clouds (black).



Figure 4. NO<sub>2</sub> S/N as a function of observing phase and distance for the O<sub>2</sub>-rich and O<sub>2</sub>-poor K-star scenarios. The black line marks a detectable S/N of 3. We use integration periods of 2 days after the flare and 10 days before the flare in this plot. These calculations include clouds.

competition between the  $NO<sub>2</sub>$  signal decaying quickly and more integration time leading to smaller noise. This is the case both for calculations including clouds and without clouds. Interestingly, the simulation with clouds yields a higher  $S/N$ for the observation because it has a smaller S/N before the flare. In both cases, two days after the flare and as many days as possible before the flare are the optimal integration periods.

Figure 4 shows the  $S/N$  as a function of orbital phase and distance for both the  $O_2$ -rich and  $O_2$ -poor K-star scenarios. Interestingly, the  $S/N$  is only weakly dependent on both  $O<sub>2</sub>$  level and phase. The weak phase dependence is due to a trade off between more reflected light at smaller phases and a larger NO<sub>2</sub> column on the night side (Figure [2](#page-2-0)), which is more visible at larger phases. We find that the  $NO<sub>2</sub>$  response to a large flare will be

<span id="page-4-0"></span>observable for K-star planets within about 3–4 pc regardless of their oxygen content.

#### 4. Caveats

In this work we were limited to species that PSG includes, among which  $NO<sub>2</sub>$  was the most detectable. However, both  $NO$ and  $HNO<sub>3</sub>$  have absorption cross sections in the UV and visible at least as strong as  $NO<sub>2</sub>$  (Keller-Rudek et al. [2013](#page-5-0)), and would be useful molecules to investigate in the future.

NASA's IROUV direct-imaging space telescope is still in planning stages. We chose best-guess parameter options in PSG meant to correspond to how IROUV is currently being imagined (National Research Council [2021](#page-5-0)). Our conclusions might change if the IROUV specifications change, and our work may need to be repeated as the IROUV plan becomes more concrete.

The model data on which this paper is based assume stellar flares and proton events with fixed proton energy spectra derived from a series of specific solar observations (see also Segura et al. [2010](#page-5-0) and Tilley et al. [2019](#page-5-0)). Proton flux and highenergy photon flux dependencies will be different depending on the stellar spectral type and the precise nature of the flaring event (Herbst et al. [2019](#page-5-0); Hu et al. [2022](#page-5-0)). Significant vertical variations in the gaseous mean molecular densities and ion compositions suggest that sporadic distributions of proton energy spectra would lead to different cumulative photolytic/ photochemical effects by virtue of deeper (or shallower) depositions of the incident charged particles.

The majority of studies using terrestrial photochemical models have considered the effects of flares on  $HO_x$ ,  $NO_x$ , NO<sub>y</sub>, and O<sub>3</sub>, or species in N<sub>2</sub>-O<sub>2</sub>-H<sub>2</sub>-CO<sub>2</sub>-rich atmospheres. As future instruments will be observing planetary systems at various ages in their evolution, it is also important to perform similar assessments for flare-modulated compositions akin to early Earth by including reduced species. Such a study will necessitate using a more flexible GCM with a chemical framework capable of simulating the effects of ion chemistry on the formation of hydrocarbons and photochemical haze (Estrela & Valio 2018). It would also be instructive to simulate the observational consequences of stellar activity (including XUV irradiation and flares) on young planetary systems, as observing planetary-mass companions in the mid-infrared will be a high priority in the JWST GO program Cycles 1 and 2 (Hinkley et al. [2022;](#page-5-0) Miles et al. [2023](#page-5-0)). Vastly divergent pathways of atmospheric composition due to stochastic delivery/loss processes (Chen & Jacobson 2022) would also interface with the heightened flare frequencies and amplitudes during those epochs (Davenport et al. 2019). A study to explore a wide variety of initial planetary parameters would require the use of Monte Carlo calculations to test a range of flare rates on a range of timescales (e.g., Smith et al. [2004](#page-5-0)).

Our results do not allow us to determine whether flare-driven variations in  $NO<sub>2</sub>$  are more detectable for K-star planets than G-star planets because the K-star planets we simulated are tidally locked or because the stellar spectrum is different. More simulations would be required to determine the effect of rotation rate on the detectability of flare-driven variations in  $NO<sub>2</sub>$  for G-star planets.

Lastly (but perhaps most importantly), the precise effects of stellar flares on exoplanet atmospheres are complex and, to some degree, dependent on the particular model employed. Another GCM study using the Met Office unified model

coupled with a chemical framework (Ridgway et al. [2023](#page-5-0)) reached somewhat different conclusions from Chen et al.  $(2021)$  and Tilley et al.  $(2019)$  $(2019)$  $(2019)$ . They found that flares cause  $O_3$ enhancement, as opposed to depletion, as UV-driven production offsets  $O_3$  photolysis during the impulsive phase of the flaring events. These possibilities warrant model comparisons of observational predictions among different radiative transfer models, photochemical models, and climate models.

## 5. Conclusion

We find that the change in atmospheric  $NO<sub>2</sub>$  resulting from large flares should be detectable using NASAs future directimaging IROUV space observatory on Earth-like planets orbiting K stars within about 3–4 pc whether they are oxygen-rich ( $\times$ 1 present atmospheric level (PAL) of O<sub>2</sub>) or oxygen-poor ( $\times 10^{-3}$  PAL of O<sub>2</sub>). This is an exciting result because it offers the opportunity to test the predictions of stateof-the-art Earth-system models employed in exoplanetary contexts including Community Earth System Model, the Met Office Unified Model, and others. Our work suggests that such a test will not be possible for planets orbiting G-stars. These conclusions should be further examined by studies with improved stellar flare, coronal mass ejection, magnetospheric transport, and atmospheric chemistry-climate models.

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